

# Solution for the Quantum Physics 1 Exam from Jan. 23, 2003

1.

$$\rho = \frac{N}{L^2} = 2 \frac{\pi p_F^2}{h^2} = \frac{2\pi \cdot 2m(\epsilon_F - V)}{h^2} = \frac{4\pi m}{h^2} \left( \epsilon_F - \frac{kx^2}{2} - ky^2 \right)$$

$$N = \int_{x^2+2y^2 \leq 2\epsilon_F/k} \rho dx dy = \int_{x^2+2y^2 \leq 2\epsilon_F/k} \rho(x, u) dx \frac{du}{\sqrt{2}} \quad (u \equiv \sqrt{2}y, r \equiv \sqrt{x^2 + y^2})$$

$$\begin{aligned} N &= \frac{4\pi m}{\sqrt{2}h^2} \int_0^{\sqrt{2\epsilon_F/k}} \left( \epsilon_F - \frac{k}{2}r^2 \right) 2\pi r dr = \frac{8\pi^2 m}{\sqrt{2}h^2} \left[ \epsilon_F \frac{r^2}{2} - \frac{k}{8}r^4 \right]_0^{\sqrt{2\epsilon_F/k}} = \\ &= \frac{8\pi^2 m \epsilon_F^2}{\sqrt{2}h^2 k} \left( 1 - \frac{1}{2} \right) = \frac{m \epsilon_F^2}{\sqrt{2}h^2 k} \Rightarrow \epsilon_F = 2^{1/4} \hbar \sqrt{\frac{kN}{m}} \end{aligned}$$

2.

$$\Delta E = \left\langle \frac{\mathbf{p}^4}{8m^3 c^2} \right\rangle, \quad p_x = i \sqrt{\frac{m\hbar}{2}} \sqrt{\frac{k}{m}} (a_x^\dagger - a_x), \quad p_y = i \sqrt{\frac{m\hbar}{2}} \sqrt{\frac{2k}{m}} (a_y^\dagger - a_y)$$

$$\Delta E = \frac{\hbar^2 k}{32m^2 c^2} \left\langle [(a_x^\dagger - a_x)^4 + 2(a_y^\dagger - a_y)^4 + 2\sqrt{2}(a_x^\dagger - a_x)^2(a_y^\dagger - a_y)^2] \right\rangle \equiv \frac{\hbar^2 k}{32m^2 c^2} X$$

The non-zero contributions to the ground state:

$$\begin{aligned} X &= \langle 0, 0 | (a_x a_x^\dagger a_x a_x^\dagger + a_x a_x a_x^\dagger a_x^\dagger + 2a_y a_y^\dagger a_y a_y^\dagger + 2a_y a_y a_y^\dagger a_y^\dagger + 2\sqrt{2}a_x a_x^\dagger a_y a_y^\dagger) | 0, 0 \rangle = \\ &= 1 + 2 + 2 + 2 \cdot 2 + 2\sqrt{2} = 9 + 2\sqrt{2} \end{aligned}$$

So  $\Delta E_{0,0} = \frac{\hbar^2 k}{32m^2 c^2} (9 + 2\sqrt{2})$ . And for the first excited state:

$$\begin{aligned} X &= \langle 1, 0 | (a_x a_x^\dagger a_x a_x^\dagger + a_x^\dagger a_x a_x^\dagger a_x + a_x a_x a_x^\dagger a_x^\dagger + a_x a_x^\dagger a_x^\dagger a_x + a_x^\dagger a_x a_x a_x^\dagger + 2a_y a_y^\dagger a_y a_y^\dagger + \\ &\quad + 2a_y a_y a_y^\dagger a_y^\dagger + 2\sqrt{2}a_x a_x^\dagger a_y a_y^\dagger + 2\sqrt{2}a_x^\dagger a_x a_y a_y^\dagger) | 1, 0 \rangle = \\ &= 4 + 1 + 3 \cdot 2 + 2 + 2 + 2 + 2 \cdot 2 + 2\sqrt{2} \cdot 2 + 2\sqrt{2} = 21 + 6\sqrt{2} \end{aligned}$$

so  $\Delta E_{1,0} = \frac{\hbar^2 k}{32m^2 c^2} (21 + 6\sqrt{2})$ .

3. *Note* The plane-wave wave-function cannot be normalized to 1 on the entire  $z$ -axis. Thus we may normalize it to give the particle-density at every point.

$$\psi(z \leq 0) = a \begin{pmatrix} 1 \\ 0 \\ \frac{pc}{E+mc^2} \\ 0 \end{pmatrix} e^{ipz/\hbar} + b \begin{pmatrix} 1 \\ 0 \\ -\frac{pc}{E+mc^2} \\ 0 \end{pmatrix} e^{-ipz/\hbar}, \quad \psi(z \geq 0) = f \begin{pmatrix} 1 \\ 0 \\ \frac{p'c}{E-V_0+mc^2} \\ 0 \end{pmatrix} e^{ip'z/\hbar}$$

$$p = \frac{1}{c} \sqrt{E^2 - m^2 c^4} = \frac{\sqrt{5}}{2} mc, \quad p' = \frac{1}{c} \sqrt{(E - V_0)^2 - m^2 c^4}$$

Boundary condition at  $z = 0$ :

$$a + b = f, \frac{pc}{E + mc^2}(a - b) = \frac{p'c}{E - V_0 + mc^2}f$$

$$V_0 = \frac{1}{2}mc^2 \Rightarrow p' = 0$$

$$a + b = f, a - b = 0 \Rightarrow f = 2a, \rho(z) = \psi^\dagger \psi = 4|a|^2$$

$$4. V_0 = 2mc^2 \Rightarrow p' = i\frac{\sqrt{3}}{2}mc$$

$$a + b = f, \frac{1}{\sqrt{5}}(a - b) = \frac{i\sqrt{3}/2}{3/2 - 2 + 1}f = i\sqrt{3}f \Rightarrow a - b = i\sqrt{15}f \Rightarrow f = \frac{2a}{1 + i\sqrt{15}}$$

$$\psi = f \begin{pmatrix} 1 \\ 0 \\ i\sqrt{3} \\ 0 \end{pmatrix} e^{i(i\sqrt{3}/2)mcz/\hbar} \Rightarrow \psi^\dagger \psi = (1 + 3)|f|^2 e^{-\sqrt{3}mcz/\hbar} = |a|^2 e^{-\sqrt{3}mcz/\hbar}$$

$$5. V_0 \rightarrow \infty \Rightarrow p' \rightarrow \infty \text{ (} p' \text{ pure real):}$$

$$a + b = f, \frac{1}{\sqrt{5}}(a - b) = -f \Rightarrow f = \frac{2a}{1 - \sqrt{5}}$$

$$\psi = f \begin{pmatrix} 1 \\ 0 \\ -1 \\ 0 \end{pmatrix} e^{ip'z/\hbar} \Rightarrow \psi^\dagger \psi = 2|f|^2 = \frac{2 \cdot 4|a|^2}{(1 - \sqrt{5})^2} = \frac{8|a|^2}{(1 - \sqrt{5})^2}$$

6. The same mathematical manipulations as done in the lecture (or in Berestetskii, Lifshitz and Pitaevskii's Quantum Electrodynamics) but with  $\mathbf{p}$  replaced by  $\mathbf{p} - \frac{e}{c}\mathbf{A}$  lead to

$$\hat{O} = \frac{1}{8m^2c^2} \left( \mathbf{p} - \frac{e}{c}\mathbf{A} \right)^2.$$

7.

$$|\psi\rangle = \frac{1}{\sqrt{2}}(a_B^\dagger + ia_A^\dagger)e^{-1/2} \sum_{n=0}^{\infty} \frac{(1/\sqrt{2})^n}{\sqrt{n!}} \frac{1}{\sqrt{n!}}(a_A^\dagger + ia_B^\dagger)^n |0\rangle$$

$$\begin{aligned} P(1_A, k_B) &= |\langle 1_A, k_B | \psi \rangle|^2 = \left| \frac{i}{\sqrt{2}}e^{-1/2} \left( \frac{i}{\sqrt{2}} \right)^k \frac{1}{\sqrt{k!}} + \frac{1}{\sqrt{2}}e^{-1/2} \frac{1}{i} \left( \frac{i}{\sqrt{2}} \right)^k \frac{k}{\sqrt{k!}} \right|^2 = \\ &= e^{-1} \left( \frac{1}{2} \right)^{k+1} \frac{1}{k!} (1 - k)^2 \end{aligned}$$

$$\begin{aligned} P(1_A) &= \sum_{k_B=0}^{\infty} P(1_A, k_B) = e^{-1} \sum_{k=0}^{\infty} \left( \frac{1}{2} \right)^{k+1} \frac{1}{k!} (1 - k)^2 = e^{-1} \sum_{k=0}^{\infty} \frac{1}{2^{k+1}k!} [k(k-1) - k + 1] = \\ &= \frac{1}{2}e^{-1} \left( \sum_{k=0}^{\infty} \frac{1}{2^k k!} - \frac{1}{2} \sum_{k=1}^{\infty} \frac{1}{2^{k-1}(k-1)!} + \frac{1}{4} \sum_{k=2}^{\infty} \frac{1}{2^{k-2}(k-2)!} \right) = \\ &= e^{-1} e^{1/2} \frac{1}{2} \left( 1 - \frac{1}{2} + \frac{1}{4} \right) = \frac{3}{8}e^{-1/2}. \end{aligned}$$

And from the expression for  $P(1_A, k_B)$  it is clear that  $P(1, 1) = 0$ , thus if it is given that there is one photon at B, there cannot be a single photon at A.

8. The two final states at A and at B are uncorrelated coherent states

$$|\alpha = \frac{1+i}{\sqrt{2}}\rangle = e^{-1/2} \sum_{n=0}^{\infty} \left(\frac{1+i}{\sqrt{2}}\right)^n |n\rangle,$$

$$\text{so } P(1) = \left|\langle 1|\alpha = \frac{1+i}{\sqrt{2}}\rangle\right|^2 = e^{-1} \left|\frac{1+i}{\sqrt{2}}\right|^2 = e^{-1}.$$

9. For  $\alpha_k$  and  $\alpha_k^\dagger$  to be annihilation and creation operators, respectively, we require  $[\alpha_k, \alpha_{k'}] = 0$ ,  $[\alpha_k^\dagger, \alpha_{k'}^\dagger] = 0$  and  $[\alpha_k, \alpha_{k'}^\dagger] = \delta_{kk'}$ .

$$[\alpha_k, \alpha_{k'}] = -u_k v_{k'} [a_k, a_{-k'}^\dagger] - v_k u_{k'} [a_{-k}^\dagger, a_{k'}] = (-u_k v_{-k} + v_k u_{-k}) \delta_{k,-k'} = 0 \Rightarrow u_k = u_{-k}, v_k = v_{-k}$$

$$[\alpha_k, \alpha_{k'}^\dagger] = (u_k u_{k'} - v_k v_{k'}) \delta_{kk'} = \delta_{kk'} \Rightarrow u_k^2 - v_k^2 = 1.$$

There is no vacuum for  $\alpha_k$  since  $\alpha_k |\psi\rangle = u_k a_k |\psi\rangle - v_k a_{-k}^\dagger |\psi\rangle$  and for bosons  $a_{-k}^\dagger |\psi\rangle$  never cancels.